Thermal Diffusivity Measurement of SnO$_2$-CuO Ceramic at Room Temperature

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ABSTRACT

Thermal diffusivity is a measure of rapidity of heat propagation through a material. The property is important in the understanding of gas sensor performance. Photoflash technique was used to determine thermal diffusivity of SnO$_2$-based materials with varying amount of CuO ranging from 10 to 50 mol%, at room temperature. The samples were made disc-shaped, 10 mm in diameter with thickness ranging from 2.4 to 2.8 mm. We found that the thermal diffusivity values of SnO$_2$-CuO samples ranging from $6.21 \times 10^{-2}$ to $7.51 \times 10^{-2}$ cm$^2$ s$^{-1}$ were better than the reported thermal diffusivity value of pure SnO$_2$ sample ($1.45 \times 10^{-2}$ cm$^2$ s$^{-1}$). The thermal diffusivity behaviour was supported by results from XRD and SEM.

Keywords: Photoflash technique, SnO$_2$-CuO, thermal diffusivity

INTRODUCTION

Atmospheric pollution and safety requirements for industry have led to research and development of a variety of gas sensors using different materials and technologies particularly for better sensor performance (Chang et al., 2002). An n-type semiconductor oxide such as tin dioxide (SnO$_2$) is one of the most important and extensively used materials for the detection of a number of toxic gases (Saadeddin et. al., 2006, Vasiliev et al., 1998). It is well established that the SnO$_2$-based gas sensors offer desirable attributes of cheapness, simplicity and high sensitivity (Kocemba et al., 2001). However, there are some limitations to the wide use of semiconductor gas sensors i.e. lack of selectivity, long term drift and sensitivity to air humidity. The improvement of sensor properties of these materials can be achieved by the addition of adequate dopants.

Thermal diffusivity, $\alpha$, is a measure of rapidity of the heat propagation through a material. The higher the value of thermal diffusivity of a material, the faster will be the heat propagation. For gas sensor application, it is important to first heat up the sensor to its operating temperature in order to activate the gas sensing property. For the sensor to be immediately used, we need the sensor to reach its operating temperature in a short time. So, sensors with higher thermal diffusivity material will give faster initialization time than sensors with lower thermal diffusivity. In this study, copper oxide, CuO was added to SnO$_2$ as the dopant to see its effects on the thermal diffusivity.

Nowadays, several different techniques for the determination of the thermal diffusivity may be found in the literatures. The proper method is usually selected by considering the
specimen size, shape, temperature characteristics and temperature range. The photoflash technique which has been developed by Parker et al. (1961) has been widely used due to the ease of sample preparation, less material requirements and fast measurements. One typically obtains accuracy within ± 5% when implementing careful measurement techniques.

MATERIALS AND METHODS

SnO$_2$-CuO ceramic was prepared using solid state route (Moon et al., 2001). The raw materials used were SnO$_2$ (99.9%, Alfa Aesar) and CuO (extra pure, Scharlau). The details of the samples composition are shown in Table 1.

<table>
<thead>
<tr>
<th>Sample Name</th>
<th>Composition (mol%)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>SnO$_2$</td>
</tr>
<tr>
<td>SC1</td>
<td>90</td>
</tr>
<tr>
<td>SC2</td>
<td>80</td>
</tr>
<tr>
<td>SC3</td>
<td>70</td>
</tr>
<tr>
<td>SC4</td>
<td>60</td>
</tr>
<tr>
<td>SC5</td>
<td>50</td>
</tr>
</tbody>
</table>

All the powders were weighed according to their compositions, wet mixed and stirred for 24 hours. The slurry was dried in an oven at 120°C for 10 hours. The mixtures were pre-sintered at 850°C for 2 hours at the heating and cooling rate of 3°C/min$^{-1}$. The samples were sieved using a 75 µm sieve. The green bodies were pressed into pellets of 10 mm in diameter with thickness ranging from 2.4 to 2.8 mm, at the pressure of 20 kN. Samples were then sintered at 1000°C for 3 hours.

The X-ray diffraction analysis was carried out for all samples using a Philip model 7602 EA Almelo X-ray Diffractometer (CuKα, λ=1.5418 Å) at range of 20 to 80 degree. The data was collected and analyzed with a Philips X’Pert Data Collector (Version 1.2d) software. The microstructure of the samples was observed by using a JEOL model 6400 Scanning Electron Microscope (SEM). To protect samples from charging by the electron beam during observation, the samples were coated with a thin gold film. Each sample was scanned at 10 kV accelerating voltage.

An important property of any material is its density. Measuring sample’s density is one way to investigate the porosity of the sample. The increase of density for the sample could be due to the sufficient space or vacancies for an effective transportation of metal ions and the decrease of density could indicate the increased number and size of pores trapped within the grains. All sintered samples had their dimensions measured and were dry weighed. The geometrical density calculation is defined as:

\[
D = \frac{m}{\pi \left(\frac{d}{2}\right)^2 L}
\]
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where $m$, $d$ and $L$ are the mass, diameter and thickness of the pellet sample, respectively. The theoretical density of the samples was calculated using the equation:

$$\rho_{\text{theo}} = \frac{m_s}{(\frac{V_{\text{SnO}_2}}{V_{\text{SnO}_2} + V_{\text{CuO}}})} \quad (2)$$

where $\rho_{\text{theo}}$, $m_s$, $V_{\text{SnO}_2}$ and $V_{\text{CuO}}$ are the theoretical density, the mass of the mixture of sample, the volume fraction of $\text{SnO}_2$ and the volume fraction of $\text{CuO}$, respectively.

Thermal diffusivity of $\text{SnO}_2$-$\text{CuO}$ ceramics was measured by using the photoflash technique at room temperature. The sample was polished using abrasive silicone carbide sandpaper to obtain a flat surface and also to get rid of impurities. Calibration was performed using standard Aluminium sample before starting the experiment. The schematic diagram of this technique is shown in Fig. 1. The front surface of the sample was exposed to a short burst of radiant energy from photoflash. The transient temperature curve of the rear surface of the sample was then measured with a thermocouple.

For photoflash technique, two ways of determining thermal diffusivity have been deduced by Parker et al. (1961) from the equation below.

$$V = 1 + 2 \sum_{n=1}^{\infty} (-1)^n \omega^n \exp(-n^2 \omega) \quad (3)$$

When $V$ is equal to 0.5, $\omega$ is equal to 1.37, so

$$\alpha = 1.37 \frac{L^2}{\pi t_{1/2}} \quad (4)$$

where $t_{1/2}$ is the time required for the back surface to reach half of the maximum temperature rise. Alternatively, $t_x$, the tangent to the temperature-time curve at the point of maximum gradient can be extrapolated to the time axis where the intercept, occur at $\omega = 0.48$ from which

$$\alpha = 0.48 \frac{L^2}{\pi t_x} \quad (5)$$
where $t_x$ is the time axis intercept versus time curve. Actually we may also use any percentage rise as the reference point for the calculation of thermal diffusivity as given below:

$$\alpha = k \frac{L^2}{t}$$  \hspace{1cm} (6)

where $K_x$ is a constant corresponding to an $x$ percentage rise of temperature at the rear surface of slab and $t_x$ is the elapsed time to an $x$ percentage rise.

In the calculation of thermal diffusivity, it is advisable to rely on a single point from the whole curve, as in the case with expression [4]. Expression [6] permits the use of the portion of the response curve, which has been least affected by interferences. The use of expression [6] is therefore strongly recommended.

The characteristics rise time is important in order to quantify whether the corrections for finite pulse effect should be applied to the measurement. For the finite pulse time effect considerations, the pulse duration, $\tau$ (5ms) of the photoflash was compared with the characteristic rise time, $t_r$ of the sample. The finite pulse effect in this study was always less than 1 ($\tau/t_r < 1$), so the finite pulse time correction is negligible (Maglic et al. 1992).

Heat loss usually occurs in the flash method. Radiation heat loss corrections can be done based on Clark and Taylor (1975) rise-curve data by employing the ratio technique. For the $t_{0.75}/t_{0.25}$ ratio i.e, the time to reach 75% of the maximum divided by the time to reach 25% of the maximum, the ideal value of the ratio is 2.272. This ratio must be determined from the experimental data. Then the correction factor, $K_R$ can be calculated from the equation

$$K_R = -0.3461467 + 0.361578 \left( \frac{t_{0.75}}{t_{0.25}} \right) - 0.06520543 \left( \frac{t_{0.75}}{t_{0.25}} \right)^2$$  \hspace{1cm} (7)

The corrected value of thermal diffusivity at the half time will be

$$\alpha_{corrected} = \frac{\alpha_{0.5} K_R}{0.13885}$$  \hspace{1cm} (8)

Besides that, corrections by using many other ratios can also be used if the constants for calculating the correction factor are provided.

**RESULTS AND DISCUSSION**

The phase composition analyzed by XRD, showed that there were no impurity phases in SnO$_2$-CuO system, *Fig. 2*. All peaks are presented by SnO$_2$ and CuO (Zhang and Liu, 2000). The absence of CuO peaks in the SCI sample may be due to the smaller X-ray scattering power of CuO than SnO$_2$. The CuO peak was observed for SC$_2$, but the intensity of CuO peak remains low, suggesting that the concentration of CuO phase in SnO$_2$–CuO is small. Strong peaks of CuO (-111, 111) were observed in SC3 (at $2\theta = 35.57^\circ$, 38.77$^\circ$, respectively). The intensity of peaks corresponding to the CuO phases
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increased slightly from SC3 to SC5 while SnO$_2$ peaks (110, 101, 211) at $\theta = 26.60^\circ$, $33.89^\circ$ and $51.80^\circ$, respectively were gradually decreased when the concentration of CuO phase in SnO$_2$-CuO were increased.

Fig. 3 shows one of the thermograms obtained from the thermal diffusivity measurements. The rest of the data obtained in this study are presented in Table 2 which show the characteristic rise time and corrected thermal diffusivity, $\alpha_c$ value at different composition of CuO mixed with SnO$_2$. Fig. 4 shows the relationship between thermal diffusivity and SnO$_2$-CuO system.

Fig. 2: XRD patterns of SnO$_2$-CuO samples

Fig. 3: Thermogram of SC1 sample (thickness, $L=0.2354$ cm)

Fig. 4: Thermal diffusivity vs. SnO$_2$-CuO system

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Thermal diffusivity of pure SnO$_2$ has been reported to be $1.45 \times 10^{-2}$ cm$^2$ s$^{-1}$ (Rosyaini, 2004). The introduction of CuO (10-50 mol %) into the ceramic system, resulted in thermal diffusivity range of 6.21 to 7.51 $\times$ $10^{-2}$ cm$^2$ s$^{-1}$. The addition of CuO up to 30 mol% increased the thermal diffusivity value from 1.45 $\times$ $10^{-2}$ cm$^2$ s$^{-1}$ (pure SnO$_2$) to 7.51 $\times$ $10^{-2}$ cm$^2$ s$^{-1}$ (SC3). However, further addition of CuO beyond 30 mol% decreases thermal diffusivity value to 6.21 $\times$ $10^{-2}$ cm$^2$ s$^{-1}$.

The influence of CuO on the sintered samples densities, theoretical densities, relative densities and porosities is shown in Table 3 while Fig. 5 shows the relative density as a function of mole percentage of CuO addition. Fig. 6 shows the porosity percentage as a function of mole percentage of CuO addition.

Based on the density measurement, the density of the sintered sample decreases gradually as CuO was increased (10-50 mol %). The highest density ($\rho = 5.87$ gcm$^{-3}$), was obtained for the SC1 sample where its relative density was about 85% of the theoretical density, while SC5 gave the lowest density ($\rho = 5.39$ gcm-3) where its relative density was about 80% of the theoretical density. This indicates that more pores were present in this sample as more CuO was introduced.

It is well understood that heat propagation will be retarded as they encounter grain boundaries. Every grain boundary is a potential site for phonon scattering. For heat
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### TABLE 3
Influence of CuO on the sintered samples densities, theoretical densities, relative densities and porosities

<table>
<thead>
<tr>
<th>Sample</th>
<th>Sintered Density (g/cm$^3$)</th>
<th>Theoretical Density (g/cm$^3$)</th>
<th>Relative Density (%)</th>
<th>Porosity (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>SC1</td>
<td>5.87</td>
<td>6.92</td>
<td>84.86</td>
<td>15.14</td>
</tr>
<tr>
<td>SC2</td>
<td>5.72</td>
<td>6.89</td>
<td>83.04</td>
<td>16.96</td>
</tr>
<tr>
<td>SC3</td>
<td>5.63</td>
<td>6.85</td>
<td>82.09</td>
<td>17.91</td>
</tr>
<tr>
<td>SC4</td>
<td>5.46</td>
<td>6.81</td>
<td>80.07</td>
<td>19.93</td>
</tr>
<tr>
<td>SC5</td>
<td>5.39</td>
<td>6.77</td>
<td>79.63</td>
<td>20.37</td>
</tr>
</tbody>
</table>

Fig. 5: Influence of CuO addition to relative density

![Graph showing the influence of CuO addition on relative density.](image)

Fig. 6: Influence of CuO addition to porosity

![Graph showing the influence of CuO addition on porosity.](image)
propagation in two samples with two different grain sizes at a same micrometer distance, the one with bigger grain size will have less grain boundaries for the heat propagation to encounter than the other sample with smaller grain size which has more grain boundaries to be encountered. Hence, phonon scattering will be less in the sample with bigger grain size; thus will show higher thermal diffusivity value. This is in agreement with the SEM micrographs in Fig. 7. Fig. 7(a), (b) and (c) represent 10 mol%, 20 mol% and 30 mol% CuO mixed with SnO₂ respectively. It shows that the particle size of these compositions was getting slightly bigger as we increased CuO from 10-30 mol% in SnO₂ samples. Therefore the addition of CuO (10-30 mol%) had contributed to the increase in the thermal diffusivity values. However, the thermal diffusivity values keep decreasing as more CuO (40-50 mol%) were doped into SnO₂ samples. The decrease in density and an increase in porosity as evidenced in SEM micrograph in Fig. 7(d) and (e), resulted in the decrease in thermal diffusivity of the sample with more than 40 mol% CuO.

Fig. 7: SEM micrograph of fractured SnO₂-CuO samples at different composition (a)10 mol% CuO (b) 20 mol% CuO (c) 30 mol% CuO (d) 40 mol% CuO (e) 50 mol% CuO, at 10,000 magnification
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**CONCLUSIONS**

The influence of CuO used as additive (10-50 mol%) on the thermal and micro-structural properties of the SnO$_2$-based ceramics were studied. The phase composition analyzed by XRD, showed that there was no impurity phases in SnO$_2$-CuO system. We found that thermal diffusivity values of SnO$_2$-CuO samples ranging from 6.21 to 7.51 $\times$ 10$^{-2}$ cm$^2$ s$^{-1}$ which is better from the reported thermal diffusivity value of pure SnO$_2$ i.e 1.45 $\times$ 10$^{-2}$ cm$^2$ s$^{-1}$. So, better performance of SnO$_2$ gas sensor in terms of thermal diffusivity, can be achieved by introducing 30 mol% CuO into SnO$_2$. The thermal diffusivity behavior was explained based on the crystallite aspect and density-porosity effect and it was supported by SEM micrographs.

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**REFERENCES**


